

VARIATIONAL PRINCIPLE FOR THE SEIBERG-WITTEN EQUATIONS

DR. CELSO MELCHIADES DORIA

ABSTRACT. Originally, the Seiberg-Witten equations were described to be dual to the Yang-Mills equation. However, there is a Variational Principle from which the SW -equation and some of their analytical properties can be studied.

1. Introduction

In november of 1994, Edward Witten gave a lecture at MIT about $N = 2$ Super-symmetric Quantum Field Theory and the ideas concerning the S -duality developed in a joint work with Seiberg in [15]. In order to please the mathematicians in the audience, he applied the new ideas to the Yang-Mills Theory to show them a new pair of 1st-order PDE, named SW -monopole equation, and conjectured that the SW -theory is dual to the Yang-Mills theory; the duality being at the quantum level, since the expectation values of dual theories are equals. In topology, this means that fixed a 4-manifold its Seiberg-Witten invariants are equal to the Donaldson invariants. After 10 years, it is believed the conjecture is true, some tour the force has been done in [11] to prove it, but in its generality it is still an open question. This new pair of equations has a simpler analytical nature than the Yang-Mills equations. Even though the open question, and the fact that the physical meaning of the Seiberg-Witten equations (SW_α -eq.) is yet to be discovered, the mathematical usefulness of the equations is rather deep and efficient to understand one of the most basic phenomenon of differential topology in four dimension, namely, the existence of non-equivalent differential smooth structures on the same underlying topological manifold. It has not been efficient enough to solve either the smooth Poincaré conjecture in dimension 4 or the 11/8-conjecture, but they have been very useful to improve the understanding of the relation between 2nd-homology classes and smooth structure, the symplectic structures on 4-manifolds and to give a construction of a large number of non-equivalent smooths structures on a compact smooth 4-manifold based on the isotopic classes of knots in S^3 [5]. Also, using the SW -theory, the Thom conjecture was proved in [9] and some results in [13] were obtained in a much easier way than the one using YM -theory. Most of the simplicity coming from the SW -theory, as compared with the YM -theory, came from the fact that the structural group in SW -theory is the abelian U_1 , wherea in YM -theory the group is the non-abelian SU_2 .

During the 80's and earlies 90's, the Yang-Mills Theory was used to define a set of smooth invariants on 4-manifolds known as Donaldson invariants [2]. The Donaldson invariants sheded new light on the theory of smooth 4-manifolds, e.g.,

Key words and phrases. connections,gauge fields,4-manifolds
MSC 58J05 , 58E50.

the existence of a non enumerable set of exotics \mathbb{R}^4 . However, the Donaldson Theory relies on a difficult analysis mostly carried out in the work of Karen Uhlenbeck [19] and also of Clifford Taubes [7]. Late in the 80's, Witten, in [21], showed that the Donaldson invariants could be described as the expectation values of a Topological Quantum Field Theory. The Yang-Mills functional was first described in 1954, by Yang and Mills [22], aiming to give a general framework in which the most basic nature's interactions would fit in. The configuration space of the Yang-Mills theory is the space of connections associated to a principal G -bundle (G =structural group) P over a 4-manifold X . The Yang-Mills functional is invariant by a conformal diffeomorphism of X and by the Gauge Group, the group of automorphism of P . The Euler-Lagrange equation of the Yang-Mills functional is known as Yang-Mills equation and written as $d^*F_A = 0$ (2^{nd} -order PDE). Whenever the theory is on a principal SU_2 -bundle, the stable critical points of the Yang-Mills functional are named *Instantons* and satisfies the anti-self dual equations (*asd*) $F_A^+ = 0$, a 1^{st} -order PDE. Thanks to the conformal invariance of the *YM*-functional, any instanton in \mathbb{R}^4 can be lifted to a solution on S^4 . Once P is fixed, the set of instantons in S^4 is a smooth manifold of dimension $d = 8k - 3$, where $k = c_2(P)$ is the 2^{nd} -Chern class of P . The *asd*-equation has some analogy with the Cauchy-Riemann equations for the Harmonic equation.

As mentioned before, the Seiberg-Witten equations came out of a duality principle and not of a variational one. However, in [20] Witten introduced a functional without making any use of it in order to derive the equations or understand their analytical properties. It turns out that the stable critical points of the this functional are the solutions of the *SW*-monopole equation; it is an open question to find on a smooth manifold a sufficient condition to the existence of a *SW*-monopole.

An outstanding difference between the Yang-Mills and *SW*-functional is the fact that the *SW*-functional is not invariant by conformal diffeomorphism of the 4-manifold and the only solution in R^4 is the trivial one $(0, 0)$, as proved in 5.1.

The *SW*-equations and the *YM*-equation became elliptic by considering them moduli the gauge invariance. In fact, the analytical properties depend on a special slice of the action defined in 3.15.

The variational setting of the Seiberg-Witten equation were explored in [8] and [4]. The main purpose of these notes is to introduce and describe the Variational Principle and some of the analytical consequences to the *SW*-equations. A interesting feature for the Variational setting would be to give an interpretation, in terms of the geometry of the configuration space, to the *SW*-invariants. The article is organised as follows;

- : section 1 - Introduction
- : section 2 - Background
- : section 3 - Variational Principle
- : section 4 - Main Estimate
- : section 5 - \mathcal{H} -Condition and Palais-Smale Condition
- : section 6 - Homotopy Type of the Configuration Space
- : section 7 - Dirchlet and Neuman Problems associated to the *SW*-equation.

It is a great pleasure to contribute to the volume in honour of the 70th-birthday of Professor Djairo Figueiredo, whose contribution to the Brazilian mathematical community has been outstanding.

2. Background

2.1. **Spin^c Structure.** The space of *Spin^c* structures on X is identified with

$$Spin^c(X) = \{\alpha + \beta \in H^2(X, \mathbb{Z}) \oplus H^1(X, \mathbb{Z}_2) \mid w_2(X) = \alpha \pmod{2}\}.$$

For each $\alpha \in Spin^c(X)$, there is a representation $\rho_\alpha : SO_4 \rightarrow \mathbb{C}l_4$, induced by a *Spin^c* representation, and a pair of vector bundles $(\mathcal{S}_\alpha^+, \mathcal{L}_\alpha)$ over X (see [10]). Let P_{SO_4} be the frame bundle of X , so

- $\mathcal{S}_\alpha = P_{SO_4} \times_{\rho_\alpha} V = \mathcal{S}_\alpha^+ \oplus \mathcal{S}_\alpha^-$.
The bundle \mathcal{S}_α^+ is the positive complex spinors bundle (fibers are *Spin₄^c* – *modules* isomorphic to \mathbb{C}^2)
- $\mathcal{L}_\alpha = P_{SO_4} \times_{det(\alpha)} \mathbb{C}$.
It is called the *determinant line bundle* associated to the *Spin^c*-structure α .
($c_1(\mathcal{L}_\alpha) = \alpha$)

Thus, for each $\alpha \in Spin^c(X)$ we associate a pair of bundles

$$\alpha \in Spin^c(X) \rightsquigarrow (\mathcal{L}_\alpha, \mathcal{S}_\alpha^+).$$

From now on, we considered on X a Riemannian metric g and on \mathcal{S}_α an hermitian structure h .

Let P_α be the U_1 -principal bundle over X obtained as the frame bundle of \mathcal{L}_α ($c_1(P_\alpha) = \alpha$). Also, we consider the adjoint bundles

$$Ad(U_1) = P_{U_1} \times_{Ad} U_1 \quad ad(\mathfrak{u}_1) = P_{U_1} \times_{ad} \mathfrak{u}_1,$$

where $Ad(U_1)$ is a bundle with fiber U_1 , and $ad(\mathfrak{u}_1)$ is a vector bundle with fiber isomorphic to the Lie Algebra \mathfrak{u}_1 .

Let \mathcal{A}_α be (formally) the space of connections (covariant derivative) on \mathcal{L}_α , $\Gamma(\mathcal{S}_\alpha^+)$ the space of sections of \mathcal{S}_α^+ and $\mathcal{G}_\alpha = \Gamma(Ad(U_1))$ the gauge group acting on $\mathcal{A}_\alpha \times \Gamma(\mathcal{S}_\alpha^+)$ as follows:

$$(2.1) \quad g.(A, \phi) = (A + g^{-1}dg, g^{-1}\phi).$$

\mathcal{A}_α is an affine space which vector space structure, after fixing an origin, is isomorphic to the space $\Omega^1(ad(\mathfrak{u}_1))$ of $ad(\mathfrak{u}_1)$ -valued 1-forms. Once a connection $\nabla^0 \in \mathcal{A}_\alpha$ is fixed, a bijection $\mathcal{A}_\alpha \leftrightarrow \Omega^1(ad(\mathfrak{u}_1))$ is explicated by $\nabla^A \leftrightarrow A$, where $\nabla^A = \nabla^0 + A$. $\mathcal{G}_\alpha = Map(X, U_1)$, since $Ad(U_1) \simeq X \times U_1$. The curvature of a 1-connection form $A \in \Omega^1(ad(\mathfrak{u}_1))$ is the 2-form $F_A = dA \in \Omega^2(ad(\mathfrak{u}_1))$.

2.2. **Seiberg-Witten Monopole Equation.** Since we are in dimension 4, the vector bundle $\Omega^2(ad(\mathfrak{u}_1))$ splits as

$$(2.2) \quad \Omega_+^2(ad(\mathfrak{u}_1)) \oplus \Omega_-^2(ad(\mathfrak{u}_1)),$$

where (+) is the self-dual component and (-) the anti-self-dual.

The 1st-order *SW*-monopole equations are defined over the configuration space $\mathcal{C}_\alpha = \mathcal{A}_\alpha \times \Gamma(\mathcal{S}_\alpha^+)$ as

$$(2.3) \quad \begin{cases} D_A^+(\phi) = 0, \\ F_A^+ = \sigma(\phi), \end{cases}$$

where

- D_A^+ is the *Spin*^c-Dirac operator defined on $\Gamma(\mathcal{S}_\alpha^+)$,
- The quadratic form $\sigma : \Gamma(\mathcal{S}_\alpha^+) \rightarrow \text{End}^0(\mathcal{S}_\alpha^+)$ given by

$$(2.4) \quad \sigma(\phi) = \phi \otimes \phi^* - \frac{|\phi|^2}{2} \cdot I$$

performs the coupling of the *ASD*-equation with the *Dirac*^c operator. Locally, for $\phi = (\phi_1, \phi_2)$, the quadratic form takes the value

$$\sigma(\phi) = \begin{pmatrix} \frac{|\phi_1|^2 - |\phi_2|^2}{2} & \phi_1 \cdot \bar{\phi}_2 \\ \phi_2 \cdot \bar{\phi}_1 & \frac{|\phi_2|^2 - |\phi_1|^2}{2} \end{pmatrix}.$$

The set of solutions of (2.3), known as *SW*-monopoles space, is the space $\mathcal{F}^{-1}(0)$, where $\mathcal{F}_\alpha : \mathcal{C}_\alpha \rightarrow \Omega_+^2(X) \oplus \Gamma(\mathcal{S}_\alpha^-)$ is a map defined by

$$\mathcal{F}_\alpha(A, \phi) = (F_A^+ - \sigma(\phi), D_A^+(\phi)).$$

The SW_α -equations are \mathcal{G}_α -invariant and the map \mathcal{F} is a Fredholm map up to the gauge equivalence.

Definition 2.2.1. *A SW_α -monopole is a solution (A, ϕ) of the SW_α -monopole equation such that $\phi \neq 0$. Solution of type $(A, 0)$ comes from an anti-self-dual connection A .*

3. Variational Principle

3.1. Sobolev Spaces. As a vector bundle E over (X, g) is endowed with a metric and a covariant derivative ∇ , we define the Sobolev norm of a section $\phi \in \Omega^0(E)$ as

$$\|\phi\|_{L^{k,p}} = \sum_{|i|=0}^k \left(\int_X |\nabla^i \phi|^p \right)^{\frac{1}{p}}.$$

In this way, the $L^{k,p}$ -Sobolev Spaces of sections of E is defined as

$$L^{k,p}(E) = \{\phi \in \Omega^0(E) \mid \|\phi\|_{L^{k,p}} < \infty\}.$$

In our context, in which we fixed a connection ∇^0 on \mathcal{L}_α , a metric g on X and an hermitian structure on \mathcal{S}_α , the Sobolev Spaces on which the basic setting is made are the following;

- $\mathcal{A}_\alpha = L^{1,2}(\Omega^1(\text{ad}(\mathfrak{u}_1)))$;
- $\Gamma(\mathcal{S}_\alpha^+) = L^{1,2}(\Omega^0(X, \mathcal{S}_\alpha^+))$;
- $\mathcal{C}_\alpha = \mathcal{A}_\alpha \times \Gamma(\mathcal{S}_\alpha^+)$;
- $\mathcal{G}_\alpha = L^{2,2}(X, U_1) = L^{2,2}(\text{Map}(X, U_1))$.
(\mathcal{G}_α is an ∞ -dimensional Lie Group which Lie algebra is $\mathfrak{g} = L^{1,2}(X, \mathfrak{u}_1)$).

The Dirichlet (\mathcal{D}) and Neuman (\mathcal{N}) problems require their own configurations spaces $\mathcal{C}_\alpha^{\mathcal{D}}$ and $\mathcal{C}_\alpha^{\mathcal{N}}$, respectively. From now on, both the configuration spaces will be denoted by \mathcal{C}_α by ignoring the superscripts, unless if it needed be.

The most basic analytical results is the *Gauge Fixing Lemma* (Uhlenbeck - [19]) and the estimate 3.1, both extended by Marini, A. [12] to manifolds with boundary; it gives a clue to define a suitable slice of the \mathcal{G}_α -action.

Lemma 3.1.1. (*Gauge Fixing Lemma*) - Every connection $\hat{A} \in \mathcal{A}_\alpha$ is gauge equivalent, by a gauge transformation $g \in \mathcal{G}_\alpha$ named *Coulomb (\mathcal{C}) gauge*, to a connection $A \in \mathcal{A}_\alpha$ satisfying

- (1) $d_\tau^{*f} A_\tau = 0$ on ∂X ,
- (2) $d^* A = 0$ on X .
- (3) In the \mathcal{N} -problem, the connection A satisfies $A_\nu = 0$ ($\nu \perp \partial X$).

Corollary 3.1.2. Under the hypothesis of 3.1.1, there exists a constant $K > 0$ such that the connection A , gauge equivalent to \hat{A} by the Coulomb gauge, satisfies the following estimates:

$$(3.1) \quad \| A \|_{L^{1,p}} \leq K \cdot \| F_A \|_{L^p}$$

notation: $*_f$ is the Hodge operator in the flat metric and the index τ denotes tangential components.

3.2. Variational Formulation. The most natural functional to be considered is

$$(3.2) \quad SW(A, \phi) = \frac{1}{2} \int_X \{ | F_A^+ - \sigma(\phi) |^2 + | D_A^+(\phi) |^2 \} dv_g.$$

Clearly, the SW_α -monopoles are the stable critical points. The next set of identities are applied to expand the functional 3.2;

Proposition 3.2.1. For each $\alpha \in Spin^c(X)$, let \mathcal{L}_α be the determinant line bundle associated to α and $(A, \phi) \in \mathcal{C}_\alpha$. Also, assume that k_g =scalar curvature of (X, g) . Then,

- (1) $\langle F_A^+, \sigma(\phi) \rangle = \frac{1}{2} \langle F_A^+ \cdot \phi, \phi \rangle$.
- (2) $\langle \sigma(\phi), \sigma(\phi) \rangle = \frac{1}{4} | \phi |^4$.
- (3) *Weitzenböck formula*

$$D^2 \phi = \nabla^* \nabla \phi + \frac{k_g}{4} \phi + \frac{F_A}{2} \cdot \phi.$$

$$(4) \quad \sigma(\phi) \phi = \frac{|\phi|^2}{2} \phi.$$

- (5) The intersection form of X $Q_X : H^2(X, \mathbb{R}) \times H^2(X, \mathbb{R}) \rightarrow \mathbb{R}$ is given by

$$Q(\omega, \eta) = \int_X \omega \wedge \eta, \quad \alpha^2 = Q(\alpha, \alpha).$$

$$(6) \quad \int_X | F_A^+ |^2 dv_g = \int_X \frac{1}{2} | F_A |^2 dv_g + 2\pi^2 \alpha^2.$$

As a consequence, a new functional turns up into the scenario;

Definition 3.2.2. For each $\alpha \in \text{Spin}^c(X)$, the Seiberg-Witten Functional is the functional $\mathcal{SW}_\alpha : \mathcal{C}_\alpha \rightarrow \mathbb{R}$ given by

$$(3.3) \quad \mathcal{SW}_\alpha(A, \phi) = \int_X \left\{ \frac{1}{4} |F_A|^2 + |\nabla^A \phi|^2 + \frac{1}{8} |\phi|^4 + \frac{1}{4} k_g |\phi|^2 \right\} dv_g + \pi^2 \alpha^2,$$

where $k_g =$ scalar curvature of (X, g) .

Remark 3.2.3.

- (1) Since X is compact and $\|\phi\|_{L^4} < \|\phi\|_{L^{1,2}}$, the functional is well defined on \mathcal{C}_α ,
- (2) The \mathcal{SW}_α -functional (3.3) being Gauge invariant induces a functional $\mathcal{SW}_\alpha : \mathcal{A}_\alpha \times_{\mathcal{G}_\alpha} \Gamma(S_\alpha^+) \rightarrow \mathbb{R}$.
- (3) The \mathcal{SW}_α -functional is bounded below by 0, and it is equal to 0 if and only if either there exists a \mathcal{SW}_α -monopole or a anti-self-dual U_1 -connection.

It is an open question to find a sufficient condition to prove that there exist a \mathcal{SW}_α -monopole on a 4-manifold X . If X is symplectic, Taubes [16] proved that whenever α is the the canonical class then there is a \mathcal{SW}_α -monopole. As a consequence of the main estimate 4.0.2, there is a non-existence result for \mathcal{SW}_α -monopoles on manifolds whose scalar curvature is non-negative. A necessary condition for the existence of a \mathcal{SW}_α -monopole is the following estimate;

Proposition 3.2.4. Let $\alpha \in \text{Spin}^c(X)$ and (A, ϕ) be a \mathcal{SW}_α -monopole. So,

$$(3.4) \quad \alpha^2 \leq \frac{2}{\pi^2} v_X \cdot (k_{g,X}^-)^4$$

Proof.

$$\mathcal{SW}_\alpha(A, \phi) - \int_X k_g |\phi|^2 = \frac{1}{4} \|F_A\|_{L^2}^2 + \|\nabla^A \phi\|_{L^2}^2 + \frac{1}{8} \|\phi\|_{L^4}^4.$$

Therefore, taking $k_{g,X}^- = \min_{x \in X} k_g(x)$,

$$\mathcal{SW}_\alpha(A, \phi) + (-k_{g,X}^-) \|\phi\|_{L^2}^2 \geq \frac{1}{4} \|F_A\|_{L^2}^2 + \|\nabla^A \phi\|_{L^2}^2 + \frac{1}{8} \|\phi\|_{L^4}^4,$$

and,

$$\|\phi\|_{L^4}^4 \leq 8\mathcal{SW}_\alpha(A, \phi) + 8(k_{g,X}^-)^2 \|\phi\|_{L^2}^2,$$

where $k_{X,g}^- = \sqrt{\max\{0, -k_{X,g}^-\}}$. It follows from inequality (7.1) that

$$\|\phi\|_{L^2}^4 \leq 8v_X \mathcal{SW}_\alpha(A, \phi) + 8v_X \cdot (k_{g,X}^-)^2 \|\phi\|_{L^2}^2.$$

So,

$$(3.5) \quad \|\phi\|_{L^2}^4 - 8v_X \cdot (k_{g,X}^-)^2 \|\phi\|_{L^2}^2 - 8v_X \mathcal{SW}_\alpha(A, \phi) \leq 0.$$

Let's consider the quadratic function $f : \mathbb{R} \rightarrow \mathbb{R}$

$$f(x) = x^2 - 8v_X \cdot (k_{g,X}^-)^2 x - 8v_X \mathcal{SW}_\alpha(A, \phi).$$

If the inequality $f(x) \leq 0$ is not satisfied by any $x \in \mathbb{R}$, then we would have $\phi = 0$. The discriminant of f is

$$\Delta = 32v_X \left(2v_X \cdot (k_{g,X}^-)^4 + \mathcal{SW}_\alpha(A, \phi) \right).$$

The inequality (3.5) admits solution if and only if $\Delta \geq 0$, since $f''(x) > 0$ and $\min_{x \in \mathbb{R}} f(x) = -\frac{\Delta}{4}$. Therefore,

$$\mathcal{SW}_\alpha(A, \phi) \geq -2v_X \cdot (k_{g,X}^-)^4$$

Since $\mathcal{SW}_\alpha(A, \phi) = -\pi^2 \alpha^2$, the lower upper bound of the \mathcal{SW}_α -functional is

$$\mathcal{SW}_\alpha(A, \phi) \geq \max\{-\pi^2 \alpha^2, -2v_X \cdot (k_{g,X}^-)^4\}.$$

In this way, if (A, ϕ) is a \mathcal{SW}_α -monopole satisfying the 1st-order \mathcal{SW}_α -equation then $\mathcal{SW}_\alpha(A, \phi) = -\pi^2 \alpha^2$, where

$$\alpha^2 \leq \frac{2}{\pi^2} v_X \cdot (k_{g,X}^-)^4.$$

The L^2 -norm of a spinor field turns out to be bounded, as shown the identity below;

$$(3.6) \quad 4v_X \cdot (k_{g,X}^-)^2 - 2\sqrt{4v_X \left[2v_X (k_{g,X}^-)^4 + \mathcal{SW}_\alpha(A, \phi) \right]} \leq \|\phi\|_{L^2}^2 \leq$$

$$(3.7) \quad \leq 4v_X \cdot (k_{g,X}^-)^2 + 2\sqrt{4v_X \left[2v_X \cdot (k_{g,X}^-)^4 + \mathcal{SW}_\alpha(A, \phi) \right]}$$

□

The Euler-Lagrange equations of the \mathcal{SW}_α -functional (3.3) are

$$(3.8) \quad \Delta_A \phi + \frac{|\phi|^2}{4} \phi + \frac{k_g}{4} \phi = 0,$$

$$(3.9) \quad d^* F_A + 4\Phi^*(\nabla^A \phi) = 0,$$

where $\Phi : \Omega^1(\mathfrak{u}_1) \rightarrow \Omega^1(\mathcal{S}_\alpha^+)$. The dual operator $\Phi^* : \Omega^1(\mathcal{S}_\alpha^+) \rightarrow \Omega^1(\mathfrak{u}_1)$ is locally, in a orthonormal basis $\{\eta^i\}_{1 \leq i \leq 4}$ of T^*X , written as

$$(3.10) \quad \Phi^*(\nabla^A \phi) = \sum_{i=1}^4 \langle \nabla_i^A \phi, \phi \rangle \eta^i, \quad \text{where } \nabla_i^A = \nabla_{X_i}^A \quad (\eta_i(X_j) = \delta_{ij}).$$

The equations above are referred as the \mathcal{SW}_α -equations.

Remark 3.2.5. The \mathcal{G}_α -action on \mathcal{C}_α has the following properties;

- (1) the \mathcal{SW}_α -functional is \mathcal{G}_α -invariant.
- (2) the \mathcal{G}_α -action on \mathcal{C}_α induces on $T\mathcal{C}_\alpha$ a \mathcal{G}_α -action as follows:
let $(\Lambda, V) \in T_{(A, \phi)}\mathcal{C}_\alpha$ and $g \in \mathcal{G}_\alpha$,

$$g \cdot (\Lambda, V) = (\Lambda, g^{-1}V) \in T_{g \cdot (A, \phi)}\mathcal{C}_\alpha.$$

Consequently, $d(\mathcal{SW}_\alpha)_{g \cdot (A, \phi)}(g \cdot (\Lambda, V)) = d(\mathcal{SW}_\alpha)_{(A, \phi)}(\Lambda, V)$.

The tangent bundle TC_α decomposes as

$$TC_\alpha = \Omega^1(ad(\mathbf{u}_1)) \oplus \Gamma(\mathcal{S}_\alpha^+).$$

In this way, the 1-form $d\mathcal{S}\mathcal{W}_\alpha \in \Omega^1(\mathcal{C}_\alpha)$ can be decomposed as $d\mathcal{S}\mathcal{W}_\alpha = d_1\mathcal{S}\mathcal{W}_\alpha + d_2\mathcal{S}\mathcal{W}_\alpha$, where

$$d_1(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} : \Omega^1(ad(\mathbf{u}_1)) \rightarrow \mathbb{R}, \quad d_1(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} \cdot \Lambda = d(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} \cdot (\Lambda, 0)$$

$$d_2(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} : \Gamma(\mathcal{S}_\alpha^+) \rightarrow \mathbb{R}, \quad d_2(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} \cdot V = d(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} \cdot (0, V).$$

By performing the computations, we get

(1) for every $\Lambda \in \mathcal{A}_\alpha$,

$$(3.11) \quad d_1(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} \cdot \Lambda = \frac{1}{4} \int_X \operatorname{Re}\{\langle F_A, d_A \Lambda \rangle + 4 \langle \nabla^A(\phi), \Phi(\Lambda) \rangle\} dx,$$

where $\Phi : \Omega^1(\mathbf{u}_1) \rightarrow \Omega^1(\mathcal{S}_\alpha^+)$ is the linear operator $\Phi(\Lambda) = \Lambda(\phi)$, which dual is defined in 3.10,

(2) for every $V \in \Gamma(\mathcal{S}_\alpha^+)$,

$$(3.12) \quad d_2(\mathcal{S}\mathcal{W}_\alpha)_{(A,\phi)} \cdot V = \int_X \operatorname{Re}\{\langle \nabla^A \phi, \nabla^A V \rangle + \langle \frac{|\phi|^2 + k_g}{4} \phi, V \rangle\} dx.$$

Therefore, by taking $\operatorname{supp}(\Lambda) \subset \operatorname{int}(X)$ and $\operatorname{supp}(V) \subset \operatorname{int}(X)$, we restrict to the interior of X , and so, the gradient of the $\mathcal{S}\mathcal{W}_\alpha$ -functional at $(A, \phi) \in \mathcal{C}_\alpha$ is

$$(3.13) \quad \operatorname{grad}(\mathcal{S}\mathcal{W}_\alpha)(A, \phi) = (d_A^* F_A + 4\Phi^*(\nabla^A \phi), \Delta_A \phi + \frac{|\phi|^2 + k_g}{4} \phi).$$

It follows from the \mathcal{G}_α -action on TC_α that

$$(3.14) \quad \operatorname{grad}(\mathcal{S}\mathcal{W}_\alpha)(g \cdot (A, \phi)) = \left(d_A^* F_A + 4\Phi^*(\nabla^A \phi), g^{-1} \cdot (\Delta_A \phi + \frac{|\phi|^2 + k_g}{4} \phi) \right).$$

3.3. Coercivity of the $\mathcal{S}\mathcal{W}_\alpha$ -Functional. An important analytical aspect of the $\mathcal{S}\mathcal{W}_\alpha$ -functional is the Coercivity Lemma proved in [8];

Lemma 3.3.1. Coercivity - *For each $(A, \phi) \in \mathcal{C}_\alpha$, there exists $g \in \mathcal{G}_\alpha$ and a constant $K_\mathfrak{e}^{(A,\phi)} > 0$, where $K_\mathfrak{e}^{(A,\phi)}$ depends on (X, g) and $\mathcal{S}\mathcal{W}_\alpha(A, \phi)$, such that*

$$\|g \cdot (A, \phi)\|_{L^{1,2}} < K_\mathfrak{e}^{(A,\phi)}.$$

Proof. lemma 2.3 in [8]. The gauge transform is the Coulomb one given in the Gauge Fixing Lemma 3.1.1. \square

Considering the gauge invariance of the $\mathcal{S}\mathcal{W}_\alpha$ -theory, and the fact that the gauge group \mathcal{G}_α is a infinite dimensional Lie Group, we can't hope to handle any analytical question in general, we need to work on a slice for the action. So forth, we restrict the problem to the space, named Coulomb Subspace,

$$(3.15) \quad \mathcal{C}_\alpha^{\mathcal{C}} = \{(A, \phi) \in \mathcal{C}_\alpha; \| (A, \phi) \|_{L^{1,2}} < K_{\mathcal{C}}^{(A, \phi)}\},$$

4. Main Estimate

In order to pursue the strong $L^{1,2}$ -convergence for the sequence $\{(A_n, \phi_n)\}_{n \in \mathbb{Z}}$, given by condition 5.0.3, next we obtain an upper bound for $\| \phi \|_{L^\infty}$, whenever (A, ϕ) is a weak solution. Due to its fundamental importance, it is named the *Main Estimate*.

Lemma 4.0.2. *Let (A, ϕ) be a solution of either \mathcal{D} or \mathcal{N} in 7.1, so*

- (1) *If $\sigma = 0$, then there exists a constant $k_{X,g}$, depending on the Riemannian metric on X , such that*

$$(4.1) \quad \| \phi \|_\infty < k_{X,g}, \quad k_{X,g} = \sqrt{\max_{x \in X} \{0, -k_g(x)\}}$$

- (2) *If $\sigma \neq 0$, then there exist constant $c_1 = c_1(X, g)$ and $c_2 = c_2(X, g)$ such that*

$$(4.2) \quad \| \phi \|_{L^p} < c_1 + c_2 \| \sigma \|_{L^{3p}}^3.$$

In particular, if $\sigma \in L^\infty$ then $\phi \in L^\infty$

Proof. Fix $r \in \mathbb{R}$ and suppose that there is a ball $B_{r^{-1}}(x_0)$, around the point $x_0 \in X$, such that

$$| \phi(x) | > r, \quad \forall x \in B_{r^{-1}}(x_0).$$

Define

$$\eta = \begin{cases} \left(1 - \frac{r}{|\phi|}\right) \phi, & \text{if } x \in B_{r^{-1}}(x_0), \\ 0, & \text{if } x \in X - B_{r^{-1}}(x_0) \end{cases}$$

So,

$$(4.3) \quad | \eta | \leq | \phi |$$

$$\nabla \eta = r \frac{\langle \phi, \nabla \phi \rangle}{|\phi|^3} \phi + \left(1 - \frac{r}{|\phi|}\right) \nabla \phi$$

$$\Rightarrow | \nabla \eta |^2 = r^2 \frac{\langle \phi, \nabla \phi \rangle^2}{|\phi|^4} + 2r \left(1 - \frac{r}{|\phi|}\right) \frac{\langle \phi, \nabla \phi \rangle^2}{|\phi|^3} + \left(1 - \frac{r}{|\phi|}\right)^2 | \nabla \phi |^2$$

$$\Rightarrow | \nabla \eta |^2 < r^2 \frac{|\nabla \phi|^2}{|\phi|^2} + 2r \left(1 - \frac{r}{|\phi|}\right) \frac{|\nabla \phi|^2}{|\phi|} + \left(1 - \frac{r}{|\phi|}\right)^2 | \nabla \phi |^2.$$

Since $r < | \phi |$,

$$(4.4) \quad | \nabla \eta |^2 < 4 | \nabla \phi |^2.$$

Hence, by 4.3 and 4.4, $\eta \in L^{1,2}$.

The directional derivative of \mathcal{SW}_α in direction η is given by

$$d(\mathcal{SW}_\alpha)_{(A,\phi)}(0, \eta) = \int_X [\langle \nabla^A \phi, \nabla^A \eta \rangle + \frac{|\phi|^2 + k_g}{4} |\phi| (|\phi| - r)].$$

By 3.12),

$$\int_X [\langle \nabla^A \phi, \nabla^A \eta \rangle + \frac{|\phi|^2 + k_g}{4} |\phi| (|\phi| - r)] = \int_X \langle \sigma, (1 - \frac{r}{|\phi|}) \phi \rangle.$$

However,

$$\int_X \langle \nabla^A \phi, \nabla^A \eta \rangle = \int_X [r \frac{\langle \phi, \nabla^A \phi \rangle^2}{|\phi|^3} + (1 - \frac{r}{|\phi|}) |\nabla \phi|^2] > 0.$$

So,

$$\int_X \frac{|\phi|^2 + k_g}{4} |\phi| (|\phi| - r) < \int_X \langle \sigma, (1 - \frac{r}{|\phi|}) \phi \rangle < \int_X |\sigma| (|\phi| - r).$$

Hence,

$$\int_X (|\phi| - r) \left(\frac{|\phi|^2 + k_g}{4} |\phi| - |\sigma| \right) < 0.$$

Since $r < |\phi(x)|$, whenever $x \in B_{r-1}(x_0)$, it follows that

$$(4.5) \quad (|\phi|^2 + k_g) |\phi| < 4 |\sigma|, \quad \text{almost everywhere in } B_{r-1}(x_0).$$

There are two cases to be analysed independently;

$$(1) \quad \sigma = 0.$$

In this case, we get

$$(4.6) \quad (|\phi|^2 + k_g) |\phi| < 0, \quad \text{almost everywhere.}$$

The scalar curvature plays a central role here: if $k_g \geq 0$ then $\phi = 0$; otherwise,

$$|\phi| \leq \max\{0, (-k_g)^{1/2}\}.$$

Since X is compact, we let $k_{X,g} = \sqrt{\max_{x \in X} \{0, -k_g(x)\}}$, and so,

$$\|\phi\|_\infty < k_{X,g}.$$

$$(2) \quad \text{Let } \sigma \neq 0.$$

The inequality 4.5 implies that

$$|\phi|^3 + k_g |\phi| - 4 |\sigma| < 0 \quad \text{a.e.}$$

Consider the polynomial

$$Q_{\sigma(x)}(w) = w^3 + k_g w - 4 |\sigma(x)|.$$

A estimate for $|\phi|$ is obtained by estimating the largest real number w satisfying $Q_{\sigma(x)}(w) < 0$. $Q_{\sigma(x)}$ being monic implies that $\lim_{w \rightarrow \infty} Q_{\sigma(x)}(w) = +\infty$. So, either $Q_{\sigma(x)} > 0$, whenever $w > 0$, or there exist a root $\rho \in (0, \infty)$. The first case would imply that

$$Q_{\sigma(x)}(|\phi(x)|) > 0, \quad a.e.,$$

contradicting 4.5. By the same argument, there exists a root $\rho \in (0, \infty)$ such that $Q_{\sigma(x)}(w)$ changes its sign in a neighborhood of ρ . Let ρ be the largest root in $(0, \infty)$ with this property. There exist constants $c_1 = c_1(X, g)$ and c_2 such that

$$|\rho| < c_1 + c_2 |\sigma(x)|^3.$$

Consequently,

$$(4.7) \quad |\phi(x)| < c_1 + c_2 |\sigma(x)|^3, \quad a.e. \text{ in } B_{r-1}(x_0)$$

and

$$(4.8) \quad \|\phi\|_{L^p} < C_1 + C_2 \|\sigma\|_{L^{3p}}^3, \quad \text{restricted to } B_{r-1}(x_0)$$

where C_1, C_2 are constants depending on $vol(B_{r-1}(x_0))$.

The inequality 4.8 can be extended over X by using a C^∞ partition of unity. Moreover, if $\sigma \in L^\infty$, then

$$(4.9) \quad \|\phi\|_\infty < C_1 + C_2 \|\sigma\|_\infty^3,$$

where C_1, C_2 are constants depending on $vol(X)$.

□

5. \mathcal{H} -Condition and Palais-Smale Condition

In the variational formulation, the problems \mathcal{D} and \mathcal{N} (7.1) are written as

$$(5.1) \quad (\mathcal{D}) = \begin{cases} grad(\mathcal{SW}_\alpha)(A, \phi) = (\Theta, \sigma), \\ (A, \phi) |_{\partial X} \stackrel{\text{gauge}}{\sim} (A_0, \phi_0). \end{cases} \quad (\mathcal{N}) = \begin{cases} grad(\mathcal{SW}_\alpha)(A, \phi) = (\Theta, \sigma), \\ i^*(F_A) = 0, \nabla_n^A \phi = 0, \end{cases}$$

The equations in 5.1 may not admit a solution for any pair $(\Theta, \sigma) \in \Omega^1(ad(\mathbf{u}_1)) \oplus \Gamma(\mathcal{S}_\alpha^+)$. In finite dimension, if we consider a function $f : X \rightarrow \mathbb{R}$, the analogous question would be to find a point $p \in X$ such that, for a fixed vector u , $grad(f)(p) = u$. This question is more subtle if f is invariant under a Lie group action on X . Therefore, we need the hypothesis below on the pair $(\Theta, \sigma) \in \Omega^1(ad(\mathbf{u}_1)) \oplus \Gamma(\mathcal{S}_\alpha^+)$;

Condition 5.0.3. (\mathcal{H}) - Let $(\Theta, \sigma) \in L^{1,2}(\Omega^1(ad(\mathbf{u}_1))) \oplus (L^{1,2}(\Gamma(\mathcal{S}_\alpha^+)) \cap L^\infty(\Gamma(\mathcal{S}_\alpha^+)))$ be a pair such that there exists a sequence $\{(A_n, \phi_n)\}_{n \in \mathbb{Z}} \subset \mathcal{C}_\alpha^{\mathcal{E}}$ (3.15) with the following properties;

- (1) $\{(A_n, \phi_n)\}_{n \in \mathbb{Z}} \subset L^{1,2}(\mathcal{A}_\alpha) \times (L^{1,2}(\Gamma(\mathcal{S}_\alpha^+)) \cup L^\infty(\Gamma(\mathcal{S}_\alpha^+)))$ and there exists a constant $c_\infty > 0$ such that, for all $n \in \mathbb{Z}$, $\|\phi_n\|_\infty < c_\infty$.
- (2) there exists $c \in \mathbb{R}$ such that, for all $n \in \mathbb{Z}$, $\mathcal{SW}_\alpha(A_n, \phi_n) < c$,
- (3) the sequence $\{d(\mathcal{SW}_\alpha)_{(A_n, \phi_n)}\}_{n \in \mathbb{Z}} \subset (L^{1,2}(\Omega^1(ad(\mathbf{u}_1))) \oplus L^{1,2}(\Gamma(\mathcal{S}_\alpha^+)))^*$, of linear functionals, converges weakly to

$$L_\Theta + L_\sigma : TC_\alpha \rightarrow \mathbb{R},$$

where

$$L_{\Theta}(\Lambda) = \int_X \langle \Theta, \Lambda \rangle, \quad L_{\sigma}(V) = \int_X \langle \sigma, V \rangle.$$

As a consequence of 3.3.1, the sequence $\{(A_n, \phi_n)\}_{n \in \mathbb{Z}}$ given by the \mathcal{H} -condition has the following properties;

- (1) converges to a pair (A, ϕ) weakly in \mathcal{C}_{α} ,
- (2) converges to a pair (A, ϕ) weakly in $L^4(\mathcal{A}_{\alpha} \times \Gamma(\mathcal{S}_{\alpha}^+))$,
- (3) converges to a pair (A, ϕ) strongly in $L^p(\mathcal{A}_{\alpha} \times \Gamma(\mathcal{S}_{\alpha}^+))$, for every $p < 4$.
- (4) The limit $(A, \phi) \in L^2(\mathcal{A}_{\alpha} \times \Gamma(\mathcal{S}_{\alpha}^+))$, obtained as a limit of the sequence $\{(A_n, \phi_n)\}_{n \in \mathbb{Z}}$, is a weak solution of 7.1 ([4]).

It turns out that the sequence $\{(A_n, \phi_n)\}_{n \in \mathbb{Z}}$, given by the \mathcal{H} -condition 5.0.3, converges strongly. The proof in [4] use the main estimate 4.0.2 and the fact that

$$\lim_{n \rightarrow \infty} \int_X \langle \Phi^*(\nabla^{A_n} \phi_n), A_n - A \rangle = 0.$$

Theorem 5.0.4. [4] *Let (Θ, σ) be a pair satisfying the \mathcal{H} - condition 5.0.3. Then, the sequence $\{(A_n, \phi_n)\}_{n \in \mathbb{Z}}$, given by 5.0.3, converges strongly to $(A, \phi) \in \mathcal{C}_{\alpha}$.*

Corollary 5.0.5. [8] *The \mathcal{SW}_{α} -functional satisfies the Palais-Smale Condition.*

5.1 - \mathcal{SW}_{α} -Monopoles in $X = \mathbb{R}^4$

It comes out of the Main Estimate that in \mathbb{R}^4 the only solution, up to gauge equivalence, to the \mathcal{SW}_{α} -equations is $(0, 0)$. The scalar curvature of \mathbb{R}^4 being $k_g = 0$ implies, by the Main Estimate,, that any solution has type $(A, 0)$, where $d^*F_A = 0$. However, by Hodge theory $F_A = 0$, so the only solution, up to gauge equivalence, is $(0, 0)$.

6. Homotopy Type of the Configuration Space

In this section, let's consider X a boundaryless smooth 4-manifold. Considering that the \mathcal{SW}_{α} -functional satisfies the Palais-Smale condition it is natural to ask about the existence of non-stable critical points. This is achieved once we know the homotopy type of \mathcal{C}_{α} . The embedding of the Jacobian Torus

$$i : T^{b_1(X)} = \frac{H^1(X, \mathbb{R})}{H^1(X, \mathbb{Z})} \hookrightarrow \mathcal{A}_{\alpha} \times_{\mathcal{G}_{\alpha}} \Gamma(\mathcal{S}_{\alpha}^+), \quad b_1(X) = \dim_{\mathbb{R}} H^1(X, \mathbb{R}),$$

defined in 6.0.6, and the variational formulation of the \mathcal{SW}_{α} -equations together give us a interpretation to the topology of $\mathcal{A}_{\alpha} \times_{\mathcal{G}_{\alpha}} \Gamma(\mathcal{S}_{\alpha}^+)$;

Proposition 6.0.6. *Let X be a closed, smooth 4-manifold. The solutions of $d^*F_A = 0$, module the \mathcal{G}_{α} -action, define the Jacobian Torus*

$$T^{b_1(X)} = \frac{H^1(X, \mathbb{R})}{H^1(X, \mathbb{Z})}, \quad b_1(X) = \dim_{\mathbb{R}} H^1(X, \mathbb{Z}).$$

Proof. Let's recall that $\alpha = \frac{i}{2\pi} \int F_A$. The equation $d^*F_A = 0$ implies that F_A is an harmonic 2-form, and by Hodge theory, it is the only one. Let A and B be solutions and consider $B = A + b$, so,

$$d^*F_B + d^*F_A + d^*db = 0 \quad \Rightarrow \quad db = 0,$$

from where we can associate $B \rightsquigarrow b \in H^1(X, \mathbb{R})$ (and $F_B = F_A$).

If a connection B_1 is gauge equivalent to B_2 , then there exists $g \in \mathcal{G}_\alpha$ such that $B = A + g^{-1}dg$ and $F_B = F_A$. However, the 1-form $g^{-1}dg \in H^1(X, \mathbb{Z})$. Consequently, if b_1, b_2 are the respective elements in $H^1(X, \mathbb{R})$, then $b_2 = b_1$ in $\frac{H^1(X, \mathbb{R})}{H^1(X, \mathbb{Z})}$. \square

The relation among the critical set of \mathcal{SW}_α -functional and the homotopy of $\mathcal{A}_\alpha \times_{\mathcal{G}_\alpha} \Gamma(S_\alpha^+)$ is described by the next result;

Theorem 6.0.7. [3] *Let X be a closed smooth 4-manifold endowed with a riemannian metric g which scalar curvature is k_g . Let*

- (1) *If $k_g \geq 0$, then the gradient flow of the \mathcal{SW}_α -functional defines an homotopy equivalence among $\mathcal{A}_\alpha \times_{\mathcal{G}_\alpha} \Gamma(S_\alpha^+)$ and $i(T^{b_1(X)})$.*
- (2) *If $k_g < 0$, then $\mathcal{A}_\alpha \times_{\mathcal{G}_\alpha} \Gamma(S_\alpha^+)$ has the same homotopy type of $T^{b_1(X)}$.*

$(A, 0)$ is a solution of the Monopole \mathcal{SW}_α -equation (minimum for \mathcal{SW}_α) whenever $F_A^+ = 0$. It is known ([2]) that if $b_2^+ > 1$, then such solutions do not exists for a dense set of the space of metrics on X . Therefore,

- (1) As a consequence of 3.2.4, only for a finite number of classes $\alpha \in Spin^c(X)$ there exists a \mathcal{SW}_α -monopole attaining the minimum.
- (2) If $\alpha \in Spin^c(X)$ is none of the classes considered in the previous item, then

$$\inf_{(A, \phi) \in \mathcal{A}_\alpha \times_{\mathcal{G}_\alpha} \Gamma(S_\alpha^+)} \mathcal{SW}_\alpha(A, \phi) > 0.$$

7. Dirichlet and Neuman Problems associated to the \mathcal{SW}_α -Equation

The Dirichlet (\mathcal{D}) and Neumann (\mathcal{N}) boundary value problems associated to the \mathcal{SW}_α -equations are the following: Let's consider $(\Theta, \sigma) \in \Omega^1(ad(u_1)) \oplus \Gamma(\mathcal{S}_\alpha^+)$ and (A_0, ϕ_0) defined on the manifold ∂X (A_0 is a connection on $\mathcal{L}_\alpha|_{\partial X}$, ϕ_0 is a section of $\Gamma(\mathcal{S}_\alpha^+|_{\partial X})$). In this way, find $(A, \phi) \in \mathcal{C}_\alpha^{\mathcal{D}}$ satisfying \mathcal{D} and $(A, \phi) \in \mathcal{C}_\alpha^{\mathcal{N}}$ satisfying \mathcal{N} , where

$$(7.1) \quad \mathcal{D} = \begin{cases} d^*F_A + 4\Phi^*(\nabla^A\phi) = \Theta, \\ \Delta_A\phi + \frac{(|\phi|^2 + k_g)}{4}\phi = \sigma, \\ (A, \phi)|_{\partial X} \stackrel{\text{gauge}}{\sim} (A_0, \phi_0), \end{cases} \quad \mathcal{N} = \begin{cases} d^*F_A + 4\Phi^*(\nabla^A\phi) = \Theta, \\ \Delta_A\phi + \frac{(|\phi|^2 + k_g)}{4}\phi = \sigma, \\ i^*(F_A) = 0, \nabla_\nu^A\phi = 0, \end{cases}$$

where $i^*(F_A) = F_4$, $F_4 = (F_{14}, F_{24}, F_{34}, 0)$ is, locally, the 4th-component (normal to ∂X) of the 2-form of curvature in the local chart (x, U) of X ;

$x(U) = \{x = (x_1, x_2, x_3, x_4) \in \mathbb{R}^4; \|x\| < \epsilon, x_4 \geq 0\}$, and

$x(U \cap \partial X) \subset \{x \in x(U) \mid x_4 = 0\}$. Let $\{e_1, e_2, e_3, e_4\}$ be the canonical base of \mathbb{R}^4 , so $\nu = -e_4$ is the normal vector field along ∂X .

The existence of a strong solution follows from 5.0.4. Basically, the regularity follows from the corollary 3.1.2 and from the Main Estimate 4.0.2, as proved in [4];

Theorem 7.0.8. [4] *If the pair $(\Theta, \sigma) \in L^{k,2} \oplus (L^{k,2} \cap L^\infty)$ satisfies the \mathcal{H} -condition 5.0.3, then the problems \mathcal{D} and \mathcal{N} admit a C^r -regular solution (A, ϕ) , whenever $2 < k$ and $r < k$.*

REFERENCES

- [1] DONALDSON, S.K. - *The Seiberg-Witten Equations and 4-Manifold Topology*, Bull.Am.Math.Soc., New Ser. **33**, n°1 (1996), 45-70.
- [2] DONALDSON, S.K.; KRONHEIMER, P. - *The Geometry of 4-Manifold*, Oxford University Press, 1991.
- [3] DORIA, CELSO M - *The Homotopy Type of the Seiberg-Witten Configuration Space*, Bull. Soc. Paranaense de Mat. **22**, n° 2, 2004. (<http://www.spm.uem.br/spmatematica/index.htm>)
- [4] DORIA, CELSO M - *Boundary Value Problems for the 2nd-order Seiberg-Witten Equations*, Journal of Boundary Value Problems, Hindawi, **1**, 2005. (<http://bvp.hindawi.com/>)
- [5] FINTUSHEL, R.; STERN, R. - *Knots, Links and Four Manifolds*, Inventiones Mathematicae **134**, n°2 (1998), 363-400.
- [6] GILBARG, D.; TRUDINGER, N.S. - *Elliptic Partial Differential Equations of Second Order*, 2nd-edition, SCSM 224, Springer-Verlag, 1983.
- [7] JAFFA, A.; TAUBES, C. - *Vortices and Monopoles*, Progress in Physics, Birkhäuser, 1980.
- [8] JOST, J.; PENG, X.; WANG, G. - *Variational Aspects of the Seiberg-Witten Functional*, Calculus of Variation **4** (1996) , 205-218.
- [9] KRONHEIMER, P.; MROWKA, T. - *The genus of Embedded Surfaces in the Projective Plane*, Math. Res. Letters **1**, 1994, 797-808.
- [10] LAWSON, H.B.; MICHELSON, M.L. - *Spin Geometry*, Princeton University Press, 1989.
- [11] FEEHAN, M.N.; LENESS, T.G. - *SO(3) monopoles, level-one Seiberg-Witten moduli spaces, and Witten's conjecture in low degrees*, Proceedings of the 1999 Georgia Topology Conference (Athens, GA). Topology Appl. **124**, 2002, no. 2, 221–326.
- [12] MARINI, A. - *Dirichlet and Neumann Boundary Value Problems for Yang-Mills Connections*, Commuc. on Pure and Applied Math, **XLV** (1992), 1015-1050.
- [13] MORGAN, J. - *The Seiberg-Witten Equations and Applications to the Topology of Smooth Four-Manifolds*, Math. Notes **44**, Princeton Press.
- [14] PALAIS, R.S. - *Foundations of Global Non-Linear Analysis*, Benjamin, inc, 1968. **22** (1990), 114-139.
- [15] SEIBERG, N.; WITTEN, E. - *Electric-Magnetic Duality, Monopole Condensation, and Confinement in N=2 SuperSymmetric Yang-Mills Theory*, Nuclear Phys. **B426** (1994).
- [16] TAUBES, C. - *GR = SW: counting curves and connections*, J. Differential Geom. **52**, 1999, no. 3, 453–609.
- [17] TAUBES, C. - *A Framework for the Morse Theory for the Yang-Mills Functional*, Inventiones Mathematicae **24**, 1988, 327-402.
- [18] TAUBES, C.; KOTSCHICK, D.; MORGAN, J.W. - *Four Manifolds without Symplectic Structure but with Non-trivial Seiberg-Witten Invariants*, Math. Res. Letter **2**, 1995, 119-124.
- [19] UHLENBECK, K. - *Connections with L^p bounds on Curvature*, Comm. Math. Phys. **83**, 1982, pp 31-42.
- [20] WITTEN, E. - *Monopoles on Four Manifolds*, Math.Res.Lett. **1**, n°6 (1994), 769-796.
- [21] WITTEN, E. - *Topological Quantum Field Theory*, Comm. Math. Phys. **117**, 1988.
- [22] - YANG, C.N.; MILLS, R.L. - *Conservation of isotopic spin and isotopic gauge invariance*, Physical Rev. (2) **96**, 1954, 191–195.

*Universidade Federal de Santa Catarina
Campus Universitário , Trindade
Florianópolis - SC , Brasil
CEP: 88.040-900
<http://www.mtm.ufsc.br>*

E-mail address: cmdoria@mtm.ufsc.br